



XV. Explicit construction of wave functions for a certain Young diagram

Two identical particles:  $\{2\}$  $\{1,1\}$



$$\hat{P}_{12} \Psi_{\{2\}} = \Psi_{\{2\}} \qquad \hat{P}_{12} \Psi_{\{1,1\}} = -\Psi_{\{1,1\}}$$

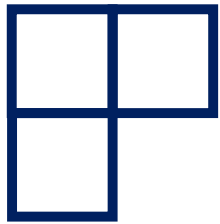
spin j and spin j yield the total spin $J = 0, 1, 2, \dots, 2j$.

$$\begin{aligned} \hat{P}_{12} |jj; JM\rangle &= \hat{P}_{12} \sum_{m_1, m_2} C_{jm_1 jm_2}^{JM} |jm_1\rangle |jm_2\rangle = \sum_{m_1, m_2} C_{jm_2 jm_1}^{JM} |jm_1\rangle |jm_2\rangle = \\ &= (-1)^{2j-J} \sum_{m_1, m_2} C_{jm_1 jm_2}^{JM} |jm_1\rangle |jm_2\rangle = (-1)^{2j-J} |jj; JM\rangle \end{aligned}$$

$\{2\}$: $J = 2j, 2j-2, 2j-4, \dots$

$\{1,1\}$: $J = 2j-1, 2j-3, \dots$

Three identical particles:  $\{3\}$  $\{1,1,1\}$

$$\hat{P}_{ij}\Psi_{\{3\}} = \Psi_{\{3\}} \quad \hat{P}_{ij}\Psi_{\{1,1,1\}} = -\Psi_{\{1,1,1\}}$$


$\{2,1\}$, an IR of S_3 with $d = 2$, transformation properties can be derived as follows:

$\psi(\sigma_1)\psi(\sigma_2)\varphi(\sigma_3)$, $\psi(\sigma_3)\psi(\sigma_2)\varphi(\sigma_1)$, $\psi(\sigma_1)\psi(\sigma_3)\varphi(\sigma_2)$ comprise a basis of a reducible representation $\{3\} + \{2,1\}$.

$$\Psi_{\{3\}}(\sigma_1, \sigma_2, \sigma_3) = [\psi(\sigma_1)\psi(\sigma_2)\varphi(\sigma_3) + \psi(\sigma_3)\psi(\sigma_2)\varphi(\sigma_1) + \psi(\sigma_1)\psi(\sigma_3)\varphi(\sigma_2)]/\sqrt{3}$$

Orthogonal states:

$$\Psi_{\{2,1\},a}(\sigma_1, \sigma_2, \sigma_3) = [\psi(\sigma_1)\psi(\sigma_2)\varphi(\sigma_3) + \psi(\sigma_3)\psi(\sigma_2)\varphi(\sigma_1) - 2\psi(\sigma_1)\psi(\sigma_3)\varphi(\sigma_2)]/\sqrt{6}$$

$$\Psi_{\{2,1\},b}(\sigma_1, \sigma_2, \sigma_3) = [\psi(\sigma_1)\psi(\sigma_2)\varphi(\sigma_3) - \psi(\sigma_3)\psi(\sigma_2)\varphi(\sigma_1)]/\sqrt{2}$$

$$\hat{P}_{ij} \begin{pmatrix} \Psi_{\{2,1\},a} \\ \Psi_{\{2,1\},b} \end{pmatrix} = \mathcal{P}_{ij} \begin{pmatrix} \Psi_{\{2,1\},a} \\ \Psi_{\{2,1\},b} \end{pmatrix}$$

$$\mathcal{P}_{12} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad \mathcal{P}_{13} = \begin{pmatrix} -\frac{1}{2} & -\frac{\sqrt{3}}{2} \\ -\frac{\sqrt{3}}{2} & \frac{1}{2} \end{pmatrix} \quad \mathcal{P}_{23} = \begin{pmatrix} -\frac{1}{2} & \frac{\sqrt{3}}{2} \\ \frac{\sqrt{3}}{2} & \frac{1}{2} \end{pmatrix}$$

Spin j and spin j couple to an intermediate momentum $K = 0, 1, \dots, 2j$.
 Then K and j couple to the final total spin $J = |K - j|, \dots, K + j - 1, K + j$.
 Although these 3 particles are identical, we write for the sake of clarity their spins as j_1, j_2, j_3 .

The basis functions (those obtained by **permutations** of 1,2,3 are equivalent) for each K :

$$|j_1, j_2(K); j_3; JM\rangle.$$

Definition of 6j-symbols:

$$\begin{aligned} \langle j_1 j_2 (j_{12}) j_3 j m | j_1, j_2 j_3 (j_{23}) j' m' \rangle &= \delta_{j j'} \delta_{m m'} U(j_1 j_2 j j_3; j_{12} j_{23}) = \\ &= \delta_{j j'} \delta_{m m'} (-1)^{j_1 + j_2 + j_3 + j} \sqrt{(2j_{12} + 1)(2j_{23} + 1)} \begin{Bmatrix} j_1 & j_2 & j_{12} \\ j_3 & j & j_{23} \end{Bmatrix}. \end{aligned}$$

therefore

$$\begin{aligned} \langle j_1 j_2 (j_{12}) j_3 j m | j_1 j_3 (j_{13}) j_2 j' m' \rangle &= \delta_{j j'} \delta_{m m'} (-1)^{j + j_1 - j_{12} - j_{13}} U(j_2 j_1 j j_3; j_{12} j_{13}) = \\ &= \delta_{j j'} \delta_{m m'} (-1)^{j_2 + j_3 + j_{12} + j_{13}} \sqrt{(2j_{12} + 1)(2j_{13} + 1)} \begin{Bmatrix} j_2 & j_1 & j_{12} \\ j_3 & j & j_{13} \end{Bmatrix}, \end{aligned}$$

Cyclic permutations of 1, 2, 3:



$$\hat{P}_{23} = \hat{P}_{12} \hat{C}, \quad \hat{P}_{13} = \hat{P}_{12} \hat{C}^{-1}$$

$$\hat{P}_{12}|j_1, j_2(K); j_3; JM\rangle = (-1)^{2j-K}|j_1, j_2(K); j_3; JM\rangle$$

$$\begin{aligned} \hat{P}_{23}|j_1, j_2(K); j_3; JM\rangle &= \hat{P}_{12}|j_2, j_3(K); j_1; JM\rangle = \\ &= (-1)^{3j+J} \hat{P}_{12} \sum_{Q=0}^{2j} \sqrt{(2K+1)(2Q+1)} \left\{ \begin{matrix} j & j & Q \\ j & J & K \end{matrix} \right\} |j_1 \cdot j_2(Q); j_3; JM\rangle = \\ &= \sum_{Q=0}^{2j} (-1)^{j+J-Q} \sqrt{(2K+1)(2Q+1)} \left\{ \begin{matrix} j & j & Q \\ j & J & K \end{matrix} \right\} |j_1 \cdot j_2(Q); j_3; JM\rangle \end{aligned}$$

$$\begin{aligned} \hat{P}_{13}|j_1 \cdot j_2(K); j_3; JM\rangle &= \\ &= \sum_{Q=0}^{2j} (-1)^{2j+K+Q} \sqrt{(2K+1)(2Q+1)} \left\{ \begin{matrix} j & j & Q \\ j & J & K \end{matrix} \right\} |j_1 \cdot j_2(Q); j_3; JM\rangle \end{aligned}$$

Then we construct linear superpositions of $|j, j(K); j; JM\rangle$ that are transformed under permutations according to IR $\{\lambda\}$.

Ammonia molecule NH_3

Total-spin states of 3 hydrogen atoms

j	J	K	$\{\lambda\}$
0	0	0	{3}
$\frac{1}{2}$	$\frac{3}{2}$	1	{3}
$\frac{1}{2}$	$\frac{1}{2}$	1	{2, 1}
$\frac{1}{2}$	$\frac{1}{2}$	0	{2, 1}

Fully deuterated ammonia ND₃

Total-spin states of 3 deuterium atoms

j	J	$\sum \alpha_K (K)\rangle$	$\{\lambda\}$
1	3	$ (2)\rangle$	$\{3\}$
1	2	$ (2)\rangle$	$\{2, 1\}$
1	2	$ (1)\rangle$	$\{2, 1\}$
1	1	$\frac{\sqrt{5}}{3} (0)\rangle + \frac{2}{3} (2)\rangle$	$\{3\}$
1	1	$\frac{2}{3} (0)\rangle - \frac{\sqrt{5}}{3} (2)\rangle$	$\{2, 1\}$
1	1	$ (1)\rangle$	$\{2, 1\}$
1	0	$ (1)\rangle$	$\{1, 1, 1\}$